Taylor classical mechanics solution manual

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Share this document with a friendTranscription:Classical Mechanics [Taylor, J.R.]Solution ManualWritten byJGSKLast UpdatedApril 4, 2018Contents1 Newton Laws of Motion 22 Bullets and charged particles 33 Moment and Angular Moment 44 Energy 55 Oscillations 66 Calculation of Variations 77 Equations of Variations Lagranges 108 Two- Central Force Problems 169 Mechanics in Non-inertial Structures 1910 Rotational Movement of Rigid Bodies 2011 Coupled Oscillators and Normal Modalities 3612 Nonlinear Mechanics and Chaos 4613 Hamiltonian Mechanics 4714 Collision Theory 6015 Special Relativity 6116 Continuous Mechanics 621 Chapter 1New Laws of Motion of Tons1.1. It is not planned to do this chapter yet.2Chapter 2Projectiles and charged particles2.1. It is not planned to do this chapter yet.3Chapter 3Moment and Angular Momentum3.1. It is not planned to do this chapter yet.4Chapter 4Energy4.1. It is not planned to do this chapter yet.5 Chapter 5Oscillations5.1. This chapter is not planned yet.6Chapter 6Calculation of Variations6.1. The shortest path between two points on a curved surface, such as the surface of a sphere, is called the ageodesic. To get a geodetic, you must first set an integral that gives the length of a path on the surface in question. This will always be similar to integral (6.2), but can be more complicated (depending on the nature of the surface) and may involve coordinates other than x and y. To illustrate this, use the sphericopolar coordinates (r , ,) to show that the length of a path joining two points on a sphere of radius R ÅL = R 211 + sin2 () 2dif (1,1) and (2,2) specify the two points and assume the path is expressed as = ().Solution. Credit: mathur/5300/5300ass2sol.pdf 6.2. Do the same as Problem 6.1 but nd the length L of a path on a cylinder of radius R, using the cylindrical-polar coordinates (,, z). Suppose the path is specified in the form = (z).Solution. Credit: mathur/5300/5300ass2sol.pdf 6.3. Consider a ray of light travelling in a vacuum from point P1 to point P2 through point Q on a flat mirror.Show that the Fermats principle implies that, in the actual path followed, Q is on the same vertical plane as P1 and P2 and obeys the law of reection, which is 1 = 2.Solution. Let us consider the luminous ray running through the path P1QP2 and use the notes given in the hints: P1 = (0,y1, 0), P2 = (x2, y2, 0) and Q = (x , 0, z) (refer to Fig. 6.8 in the manual). Then, according to Fermats's principle, light travels between two points along the path that takes less time. As the speed of light is 7 mathur/5300/5300ass2sol.pdf mathur/5300/5300ass2sol.pdfClassical Mechanics J.R.]Continuous manual solution over the entire path, minimizing travel time equals minimizing the path length L of P1QP2:L (z) =x2 + y21 + z2 + (x2 x) 2 + y22 + z2 (6.1) dLdz =zx2 + y 21 + z2+z (x2 x) 2 + y22 + z2!= 0 (6.2) Clearly the derivative is 0 only when z z z 0, that is, Q is in the xy plane (show). Following this conclusion, we now consider the light beam tied in the xy plane. The minimisation of L(z = 0) compared to x (due to the principle of Fermats) then gives:0!= dLdx =x2 + y21 (x2 x)(x2 x)2 + y22(6.3)x2 + y21=(x2 x)(x2 x)2 + y22(6.4)sin1 = sin2 (6.5) and then 1 = 2 (shown). 6.4. A light beam travels from point P1 to a refractive index mean n1 to P2 in an index medium n2, within point Q on the air interface between the two media. Showing that the principle of Fermats implies that, on the real path followed, Q is located in the same vertical plane of P1 and P2 and obeys the law Snells, thatn1 sin1 = n2 sin2. Solution. Credit: 6.5. The principle of Fermats is often referred to as the travel time of a radius of light, passing from point A to B, is minimal along the real path. Strictly speaking, it should say that time is still, not minimal. In fact, you can build situations for which time is maximum along the real path. Here's one: Consider the concave, hemisphere mirror shown in Figure 6.10 (in textbook), with A and B at opposite ends of adiameter. Consider a radius of light that travels in a vacuum from A to B with a reection to P, in the same vertical plane of A and B. According to the law of reection, the real path goes via Po point at the bottom of the hemisphere (= 0). Find the time to travel along the APB path as a function of and show that it is very good at P = Po.Solution. Credit: 8 Classical Mechanical [Taylor, J.R.] Solution Manual6.7. Consider a right circular cylinder of the R radius centered on the axis z. Find the equation giving as a z function for the geodeic (shorter path) on the cylinder between two points with cylindrical polar coordinates (R,1, z1) and (R,2, z2). Describe geodetics. Is it unique? By imagining the surface of the unmounted cylinder and tuning, he explains why geodetics has the shape it does. Solution. Credit: mathur/5300/5300ass2sol.pdf 6.8. Make sure that the speed of the roller coaster car in Example 6.2 (in textbook) is2y. (Assome wheels have unreadable mass and neglect friction.) Solution. Applying energy conservation when the car is released at height y (only potential energy) and when it is at ground level (y = 0; only kinetic energy) gives:12mv2 =my v =2yche is what we need to show. [The remaining solutions in Chapter 6 can be found at]9 mathur/5300ass2sol.pdf 7Lagranges Equations7.1. Write the Lagrangiana bullet (subject to no air resistance) inCartesian coordinates (x ,y,z), with z measured vertically upwards. Find the three Lagrange equations and prove that they are exactly what you would expect for the equations of motion. Solution. The Lagrangian can be easily identified and written as such:L = T U=12m (x2 + y2 + z2) mz (7.1) A direct application of the Euler-Lagrange equations produces:Lx=ddtL x m x = 0 (7.2) Ly=ddt L y m y = 0 (7.3) Lz=ddt L z m z = m (7.4) All three of which are intended for a free-fall bullet. Notice the negative sign on the z-direction is due to the chosen sign convention. 7.8. (a) Write the Lagrangian L (x1,x2, x1, x2) for two particles of equal mass, m1 = m2 = m, connected to the x axis and connected by a spring with potential energy U = 12kx2. [Here x is the extension of the spring, x = (x1 x2 l), where l is the length of the unreinforced springs, and I assume mass 1 remains to the right of mass 2 at all times.] (b) Rewrite L in terms of the new variables X = 12 (x1 + x2) (the CM position) and x (the extension), and write the two Lagrange equations for X and x. e) Solve for X (t) and x (t) and describe the motion. 10Classic Mechanics [Taylor, J.R.]Manual solution. (a) The Lagrangian is simply: L = 12m (x12 + x22) 12k (x1 x2 l) 2 (7.5) (b) x 7.1 x 7.1 x 1.7 x 7.1 x 2.7 x 7.1 x 2.1 x 2 l), as such:X = 12 (x1 + x2) and x = x1 x2 (7.6) x1 = X + 12 xx2 = X 12 x (7.7) We can then rewrite the Lagrangian from Eq. (7.5) as: L = 12m[(X + 12 x 12 x 12 x Now we have to solve for X (t) and x (t). Let's address the easiest of the two rst:X = 0 X (t) = v0t + X0 (7.13) where we took the liberty of introducing two integration constants that depend on the initial conditions. For the other x-coordinates , we invoke the family SHM solution: x = 2kmx x (t) = A cos (2km+) (7.14) where we have introduced two more integration constants A and . The center of mass moves with constant growth, as no external force acts upon it. The extension of the spring undergoes a simple harmonic, which means only that the two masses oscillate relative to each other. 11Classical Mechanics [Taylor, J.R.]Solution Manual7.33. A soap bar (mass) rests on a frictionless rectangular plate resting on a horizontal table. At timet = 0, I start lifting one edge of the dish so that the plate rotates around the opposite edge with constant angular velocity, and the soap begins to slide downwards towards the edge. Show that the motor equation for the soap has the shape x 2x = sieve , where x is the soap awayPage 2Browse all categories to find your favorite topic Taylor J.R. 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